**Degres Journal (2024) 9(10):51 ISSN NO:0376-8163**

### **Computational Study of Field-Dependent Pinning Force for Pure MgB2 Superconductor**

Intikhab A. Ansari

*Department of General Studies, Jubail Industrial College, P. O. Box - 10099, Jubail Industrial City-31961, Saudi Arabia* 

Accepted: 4 October 2024

### **Abstract**

The logarithmic exponential regression model is proposed for the pure  $MgB<sub>2</sub>$  superconductor in this study. The predicted values of normalized pinning force as a function of the reduced field are evaluated and compared to the measured values at a wide range of temperatures. The proposed model well described the best-suited model for the estimated predicted values. The predicted values of the reduced field dependence of normalized pinning force are in good agreement with the measured values for the pure MgB2 superconductor. Some statistical error estimations namely mean bias error, mean absolute bias error and root mean square error is calculated at different temperatures from 10−30K. The large value of the coefficient of regression,  $\mathbb{R}^2$  confirms the effectiveness and significance of the proposed model. The proposed model is applicable in this study to find the accurate and precise values of the computational data.

**Keywords:** MgB2; empirical model; pinning force; statistical errors

### **1. Introduction**

In recent twenty years, several contributions are made to enhance the performance of the  $MgB<sub>2</sub>$  superconductor. The transition temperature,  $T_c$  of the MgB<sub>2</sub> is higher than the traditional Nb-based superconductors [1]. The  $MgB<sub>2</sub>$ continues to attract the interest of scientists due to its

https://doi.org/10.5281/zenodo.13910671 Published online: 9 October 2024

remarkable properties. [2]. The high trapped field applications such as high energy particle accelerators, high field magnets, and magnetic resonance imaging require the bulk material of high critical current density,  $J<sub>c</sub>$  at high magnetic fields [3–5]. The main disadvantage of the  $MgB_2$  material is its low irreversibility,  $H_{irr}$ . Therefore, special efforts are needed to enhance its

Address correspondence to E-mail: ansari\_i@rcjy.edu.sa

superconducting properties, especially the critical current density,  $J_c$  in high magnetic fields.

The flux pinning force is an important physical factor that can be measured by the product of critical current density, *J*c, and applied field, *H*. Therefore, the study of the pinning force in bulk  $MgB_2$  is significant to drive the information on the pinning mechanism [6].

Dew-Hughes study the normalized volume pinning force as a function of the applied field to analyze the behavior of the pinning mechanism in the bulk material [7]. Furthermore, Fietz and Webb proposed the one power scaling law for normalized pinning force density to recognize the nature of pinning boundaries [8]. Pinning force scaling force analysis was performed in pure  $MgB_2$  samples with a wide range of temperatures including a spark plasma sintered one [9]. The elementary pinning force of grain boundaries is calculated using the electron scattering process and the residual resistivity [10]. In superconducting polycrystalline MgB2, electrical connectivity and flux pinning strength of grain boundaries are significant elements to find the critical current density by the Matsushita and co-workers [10]. Extrinsic twodimensional flux pinning centers have been created into MgB2 superconductors using in-situ and diffusion sintering processes, with graphene-encapsulated boron powder [11]. In our previous work, I have proposed three empirical regression models including linear, exponential, and quadratic regression models that predict the transition temperature for numerous elements doped MgB2 superconductors based on the two descriptors namely electrical connectivity index and valence energy level connectivity index [12].

Various thin films and single crystals of  $MgB<sub>2</sub>$  samples were prepared by spark-plasma sintering to evaluate the sharp narrow peak in volume pinning force vs. reduced field [13]. The pinning force scaling changes from grain boundary pinning to point pinning on increasing preparation temperature and finally the sharp, low-field peak in volume pinning force  $(F_p)$  disappears. The pinning behaviors of four different thicknesses for MgB2 thin films have been discussed by Yang and coworkers [14]. These films were prepared by a hybrid physical-chemical vapor deposition technique. It was found that MgB2 films had different growth modes in different growth stages. The effect of hydrostatic pressure up to 1.2 GPa on the critical current density and the nature of the pinning mechanism in  $MgB<sub>2</sub>$  has been investigated by collective theory [15]. It was found that the pressure increases the anisotropy and reduces the coherence length that resulting in weak interaction of the vortex cores with the pinning centers. The critical current density,  $J_c$  was improved in bulk pure  $MgB_2$ samples by optimizing the sintering conditions [16]. These results reveal a strong relationship between the microstructure and the pinning performance. Thermally activated relaxation and pinning in spark-plasma sintered MgB<sub>2</sub> superconductor were discussed [17]. Vortex pinning was investigated with the help of field dependence of the pinning force density that indicates collective pinning by normal point-like defects. The field dependence of the pinning force in different, highdensity sintered samples was analyzed for MgB<sup>2</sup> samples [18]. Within the grain boundary pinning mechanism, these samples show the most aspects of the field dependency of the critical force that fit the Dew-Hughes scaling law predictions.

In the present study, the normalized pinning force density,  $F_p/F_{p,max}$  as a function of reduced field,

 $H/H_{irr}$  have been discussed with a wide range of temperatures using empirical models for pristine MgB<sub>2</sub> superconductors. The measured  $F_p/F_{p,max}$  are compared with predicted ones using some statistical tools to calculate the error estimations Mean Bias Error (MBE), Mean Absolute Bias Error (MABE), and Root Mean Square Error (RMSE) included with Coefficient

of Regression  $(R^2)$ . These error estimations were used to determine the least error at different temperatures 10, 15, 20, 25, and 30 K.

#### **2. Proposed modeling**

The ground boundary pinning prevails in the pure  $MgB_2$ material, as discussed in the previous report [19]. In the present study, the normalized pinning force density,  $f(x)$ can be described by the following equation:

$$
f(x) = A \ln x + B e^{Cx^2}
$$
 (1)

where  $f(x) = F_p/F_{p,max}$  and  $x = H/H_{irr}$  is the reduced field.  $F_{p,max}$  is the maximum pinning force density and  $H_{irr}$  is the irreversibility field line. In addition, A, B, and C are the empirical constants.

#### *2.1 Proposed empirical model*

The logarithmic exponential regression model (LERM) attempts to model the correlation between two variables by fitting the linear combination of logarithmic function and exponential function to observe the data. One variable is considered to be an explanatory variable, and the other is to be a dependent variable. LERM supports and predicts the behavior of the given data. This is the dominant type of regression analysis to study rigorously and used extensively in practical applications.

#### *2.2 Statistical error estimation*

In the present study, the empirical model LERM is proposed. This model estimates the statistical error to the wide range of temperatures. Furthermore, the high  $R<sup>2</sup>$  value demonstrates the efficacy of the proposed model. The following are descriptions of the total sum of squares (SST) and the sum of squares of regression (SSR):

Total sum of squares  $(SST) = \sum_{i=1}^{n} \left( \overline{Y_{cm}} - Y_{cm}^{i} \right)$  $\sum_{i=1}^{n} \left( \overline{Y_{cm}} - Y_{cm}^{i} \right)^{2}$  (2)

Sum of squares of regression  
(SSR) = 
$$
\sum_{i=1}^{n} (Y_{cm}^{i} - Y_{cp}^{i})^2
$$
 (3)

here,  $Y_{cm}$  denotes the measured value and  $Y_{cp}$  indicates the predicted one, while  $Y_{cm}$  is the mean of measured value as described underneath:

$$
\overline{Y_{cm}} = \frac{1}{n} \sum_{i=1}^{n} Y_{cm}^{i}
$$
\n(4)

Coefficient of Regression (R<sup>2</sup>) = 
$$
\frac{SSR}{SST}
$$
 (5)

The coefficient of regression  $(R^2)$  illustrates how effective the best-fit is all over the data, as described by Eq (5). The empirical model with the best fit is one with a large  $R^2$  value near 1.

The residuals are the variations between the observed and predicted responses. It might be considered as the fundamentals of variation that the fitted model fails to explain. Residues, on the other hand, evaluate the random errors. As a result, the random residuals support the fit of the proposed model, whereas the regular pattern displays the least-fit model. The RMSE, MBE, and MABE are some of the statistical error estimation methods used in this study as mentioned below:

$$
\text{MBE} = \frac{1}{n} \sum_{i=1}^{n} \left( Y_m^i - Y_p^i \right) \tag{6}
$$

$$
\text{MABE} = \frac{1}{n} \sum_{i=1}^{n} \left| Y_m^{\ i} - Y_p^{\ i} \right| \tag{7}
$$

RMSE = 
$$
\sqrt{\frac{1}{n} \sum_{i=1}^{n} (Y_m{}^{i} - Y_p{}^{i})^2}
$$
 (8)

The RMSE statistic is for short-term performance, whereas the MBE figure out for the long-term. In case MBE is positive, it suggests that the predicted data is over-approximated, whereas if MBE is negative, it shows that the predicted data is under-approximated. MABE denotes the degree to which the fitting model is effective. The proposed empirical model with the least MBE, MABE, and RMSE values demonstrates the bestfitted model to the given set of data with a wide range of temperatures.

#### **3. Results and discussion**

In the proposed model, I have fitted the measured values at different temperatures ranging from 10–30K for a pure MgB2 superconductor. In this work, I have estimated the statistical errors by the SPSS program. The normalized pinning force density,  $F_p/F_{p,max}$  vs. the reduced field,  $H/H_{irr}$  measured values for pure  $MgB_2$ superconductor are taken out from my previous report [19]. The detailed procedure of the synthesis method is reported there. The critical current density,  $J_c$  was calculated from the width of the hysteresis loop by using Bean's model as mentioned in Ref [19]. In the present work, Fig. 1(a)–(e) elucidates the plot of normalized pinning force density,  $F_p/F_{p,max}$  vs. the reduced field,  $H/H_{irr}$  at different temperatures 10, 15, 20, 25, and 30K. The measured values of  $F_p/F_{p,max}$  vs.  $H/H_{irr}$ plots are fitted well with predicted values with a wide range of temperatures as shown in figures 1(a)–(e). It is remarkable from all the plots that the  $F_p/F_{p,max}$  values are enhancing sharply with the increase of  $H/H_{irr}$ values thereafter, these are decreasing exponentially.

Therefore, I can accomplish that the measured values are in agreement with the predicted one for the proposed model.

The proposed model estimates some statistical errors MBE, MABE, and RMSE, among other statistical coefficients of regression, as indicated in Table 1. The evaluated coefficients of regression  $(R^2)$ , are almost similar values 0.989 for all temperatures except 10K as shown in Table 1. These values of  $\mathbb{R}^2$  are very close to 1 that demonstrating the best-fitted values from 15−30K. The least values of statistical errors MBE, MABE, and RMSE are evaluated -0.0067, 0.02656, and 0.0346, respectively at 30K. While the highest values of MBE, MABE, and RMSE are obtained -0.0153, 0.04770, and 0.0655, respectively at a single temperature 10K as shown in Table 1. The empirical constants A, B, and C are the lowest which consists of the values 0.172, 1.422, and -2.799, respectively at 10K. I have estimated the highest peak values of  $H/H_{irr}$  0.11765, 0.13046, 0.13447, 0.14567 and 0.16003 at 10K, 15K, 20K, 25K

and 30K, respectively for predicted fitted data. It is evident that these peaks are shifting towards the high reduced field values with the increase of temperature which is in good agreement with the previous report [19]. Table 1 shows that the proposed model best described the fitted values at 30K temperature as compared to other temperatures. In addition, the measured values of  $F_p/F_{p,max}$  vs.  $H/H_{irr}$  plots are in agreement with the predicted values for the proposed model at almost all the temperatures for the pure  $MgB_2$ superconductor.

## **4. Conclusions**

I have estimated the predicted values for normalized pinning force density,  $F_p/F_{p,max}$  vs. the reduced field,  $H/H_{irr}$  plot for pure MgB<sub>2</sub> superconductor between temperatures range 10–30K. These predicted values are in agreement with the measured values that best described the proposed model. Furthermore, the evaluated statistical errors MBE, MABE, and RMSE illustrate the least values which clearly shows the bestfitted values at 30K temperature rather than the other temperatures. The coefficient of regression,  $\mathbb{R}^2$  is very close to 1 at almost all temperature ranges reveals the relevant proposed model. It is obvious from all of the plots that the  $F_p/F_{p,max}$  values improve abruptly with the increase of  $H/H_{irr}$  values, and then decreases exponentially afterward.

Table 1: Some empirical constants with their statistical values from temperature range 10−30K





Figure 1(a)-(e): The normalized pinning force density,  $F_p/F_{p,max}$  vs. the reduced field,  $H/H_{irr}$  plot for pure MgB<sub>2</sub> superconductor at temperature 10-30K.

## **Acknowledgment**

IAA thanks to Department of General Studies, Jubail

Industrial College for the support of this work.

# **Conflicts of Interest**

The author declares no conflict of interest.

# **References**

- [1] J Nagamatsu, N Nakagawa, T Muranaka, Y Zenitani and J Akimitsu, *Nature* **410** 63 (2001)
- [2] C Buzea and T Yamashita *Supercond. Sci. & Technol.* **14** R115 (2001)
- [3] K Sivasubramaniam et al. *IEEE Trans. Appl. Supercond.* **19** 1656 (2009)
- [4] M Muralidhar, K Inoue, M R Koblischka, M Tomita and M Murakami *J. Alloys Comp.* **608** 102 (2014)
- [5] A Yamamoto, A Ishihara, M Tomita, and K Kishio *Appl. Phys. Lett.* **105** 032601 (2014)
- [6] E Martínez, P Mikheenko, M Martínez-López, A Millán, A Bevan, and J S Abell *Phys. Rev. B* **75** 134515 (2007)
- [7] D Dew-Hughes *Philos. Mag. Part B* **55** 459 (1987)
- [8] W A Fietz and W W Webb *Phys. Rev.* **178** 657 (1996)
- [9] M R Koblischka et al. *J Supercond. Nov. Magn.* **33** 3333 (2020)

# **Funding**

The authors received no financial support for the research, authorship, and/or publication of this article.

## **Data availability**

The raw/processed data required to reproduce these findings cannot be shared at this time due to legal or ethical reasons.

- [10] T Matsushita et al. *Supercond. Sci. Technol.* **21** 015008 (2007)
- [11] W Li et al. *ACS Appl. Mater. Interfaces* **11** 10818 (2019)
- [12] I A Ansari and C V Rao *Materials Today: Proceedings* **45** 4417 (2021)
- [13] M Koblischka et al. AIP Adv. **10** 015035 (2020)
- [14] C Yang et al. AIP Adv. **7** 035117 (2017) [15] B Shabbir et al. *Supercond. Sci.*
- *Technol.* **28** 055001 (2015)
- [16] M Muralidhar et al. *J. of Alloys and Compounds* **649** 833 (2015)
- [17] M Jirsa et al. *Supercond. Sci. Technol.* **29** 025006 (2016)
- [18] V Sandu et al. *Scientific Reports* **11** 5951 (2021)
- [19] I A Ansari et al. *Indian J Phys.* **94** 485 (2020)